SoSe 2007 5/25/07

Prof. Dr. E. Frey Dr. T. Franosch

Lehrstuhl für Statistische Physik Biologische Physik & Weiche Materie Arnold-Sommerfeld-Zentrum für Theoretische Physik Department für Physik

T II: Elektrodynamik

(Prof. E. Frey)

Problem set 6

Tutorial 6.1 Sound waves in a fluid

The macroscopic properties of a fluid are characterized in terms of a few fields, e.g., the mass density $\rho(\vec{r},t)$, the mass current density $\vec{j}(\vec{r},t)$, the fluid velocity $\vec{v}(\vec{r},t)$, and the pressure $p(\vec{r},t)$. Euler's equations specify the field equations; the first set encodes the conservation of mass and momentum,

$$\partial_t \rho + \nabla_k j_k = 0, \quad \partial_t j_k + \nabla_l \Pi_{kl} = 0.$$
(*)

The mass current density is connected to the fluid velocity by $\vec{j}(\vec{r},t) = \rho(\vec{r},t)\vec{v}(\vec{r},t)$, and Π_{kl} denotes the momentum current tensor

$$\Pi_{kl} = \rho v_k v_l - \sigma_{kl} = \rho v_k v_l + p \delta_{kl} \,, \tag{**}$$

which closes the equations. The term $\rho v_k v_l$ is the contribution to momentum current by the inertia of the flow (the terms responsible for turbulence). The quantity $\sigma_{kl} = -p\delta_{kl} + \sigma'_{kl}$ is known as stress tensor, p denotes the pressure; last σ'_{kl} encompasses the (bulk and shear) viscous forces, i.e. dissipative processes, which are neglected in Euler's equations, $\sigma'_{kl} = 0$.

a) Demonstrate that

ĥ

$$p(\vec{r},t) = \rho_0 = const, \quad \vec{v}(\vec{r},t) = 0, \text{ and } p(\vec{r},t) = p_0 = const$$

constitutes a solution of the field equations. Show that the linearized field equations for small perturbations $\delta \rho = \rho - \rho_0$, \vec{v} , and $\delta p = p - p_0$ to this reference state read

$$\partial_t \delta \rho + \rho_0 \nabla_k v_k = 0, \qquad \rho_0 \partial_t v_k = -\nabla_k \delta p.$$

Introduce the *isothermal compressibility* κ_T that reflects the pressure increase due to compression at constant temperature to linear order, $\delta p = \delta \rho / \rho_0 \kappa_T$.

b) Derive a local conservation law, $\partial_t u + \operatorname{div} \vec{S} = 0$, for the energy density

$$u(\vec{r},t) = \frac{\rho_0}{2}\vec{v}(\vec{r},t)^2 + \frac{A}{2}\delta\rho(\vec{r},t)^2$$

for suitably chosen A relying on the approximations introduced so far. Determine the energy current density $\vec{S}(\vec{x},t)$.

c) Show that the linearized field equations allow for monochromatic longitudinal waves in \vec{v} and scalar waves in $\delta\rho(\vec{r}, t)$.





Tutorial 6.2 Polaritons

Consider the constitutive equation of the Lorentz-Drude model,

$$\partial_t^2 \vec{P}(\vec{x},t) + \frac{1}{\tau} \partial_t \vec{P}(\vec{x},t) + \omega_0^2 \vec{P}(\vec{x},t) = \frac{\omega_p^2}{4\pi} \vec{E}(\vec{x},t) \,,$$

with the relaxation time τ , characteristic frequency ω_0 and the plasma frequency ω_p .

- a) Perform a spatio-temporal Fourier transform and determine the complex susceptibility $\chi(\omega)$, with $\vec{P}(\vec{k},\omega) = \chi(\omega)\vec{E}(\vec{k},\omega)$, as well as the dielectric function $\varepsilon(\omega) = 1 + 4\pi\chi(\omega)$.
- b) Argue that the longitudinal modes follow from the zero of the dielectric function, $\varepsilon(\omega_*) = 0$, and determine the complex frequency ω_* in the case of weak damping.
- c) Ignoring the damping, $\tau \to \infty$, determine the dispersion relation of the transverse modes.
- d) Explain without calculation, in what frequency regime the damping is most important.

Problem 6.3 Doppler effect

Consider again the ideal fluid of Tutorial 6.1.

- a) Convince yourself that a fluid moving at constant velocity $\vec{v}(\vec{r},t) = \vec{v}_0$ with constant density and pressure constitutes a valid solution of the field equations (*) and (**). Linearize the field equations around this new reference state relying again on the thermodynamic relation $\delta p = \delta \rho / \rho_0 \kappa_T$.
- b) Demonstrate the existence of monochromatic plane density waves and discuss the dispersion relation $\omega = \omega(\vec{k})$ with respect to the direction of the wave propagation \vec{k} to the fluid velocity \vec{v}_0 and interpret your result.

Problem 6.4 Elastic waves

Small deformations of an isotropic elastic medium are described in terms of the vector field of displacements $\vec{u}(\vec{x},t)$, the velocities $\vec{v}(\vec{x},t)$, and the symmetric stress tensor field $\sigma_{ij}(\vec{x},t)$. The linearized field equations read

$$\partial_t u_i(\vec{x}, t) = v_i(\vec{x}, t)$$
 and $\varrho_0 \partial_t v_i(\vec{x}, t) = \nabla_j \sigma_{ij}(\vec{x}, t)$.

The field equations are closed with Hooke's law as constitutive equation,

$$\sigma_{ij} = K \delta_{ij} \operatorname{div} u + \mu \left(\nabla_i u_j + \nabla_j u_i - \frac{2}{3} \delta_{ij} \operatorname{div} u \right).$$

Here $\rho_0 > 0$ is the mass density, $\mu > 0$ denotes the shear modulus and K > 0 the bulk modulus; the inverse of K gives the compressibility.

- a) Perform a spatio-temporal Fourier transform and show that the medium supports longitudinal as well as transverse waves. Derive the corresponding dispersion relations and compare the different sound velocities.
- b) Argue that in the limit of zero shear modulus, $\mu = 0$, one recovers the hydrodynamics of an ideal fluid introduced in Tutorial 6.1.

Problem 6.5 Debye-Hückel and Thomas-Fermi theory

As a natural extension of Drude's theory of metals, consider the constitutive equation

$$\partial_t \vec{j}^{(\text{ind})}(\vec{r},t) + \frac{1}{\tau} \vec{j}^{(\text{ind})}(\vec{r},t) + c_0^2 \operatorname{grad} \rho^{(\text{ind})}(\vec{r},t) = \frac{\omega_p^2}{4\pi} \vec{E}(\vec{r},t) \,.$$

Here the new term $c_0^2 \operatorname{grad} \rho^{(\operatorname{ind})}(\vec{r},t)$ describes a restoring force similar to Euler's theory of fluids. The material is characterized by the velocity c_0 , the relaxation time τ , and the plasma frequency ω_p . Derived quantities are the screening length $\lambda_0 = c_0/\omega_p$ and the conductivity $\sigma = \omega_p^2 \tau/4\pi$.

a) Formulate a continuity equation for the scalar field

$$u_{\rm M}(\vec{r},t) = \frac{2\pi}{\omega_p^2} \left[\vec{j}^{\rm (ind)}(\vec{r},t)^2 + c_0^2 \rho^{\rm (ind)}(\vec{r},t)^2 \right] \,,$$

and specify the corresponding source term. Derive a conservation law for the total energy density consisting of the energy density of matter and the electromagnetic fields.

b) Consider an external charge Q located at the origin $\vec{r}_0 = 0$ of the medium. Discuss the induced charge density $\rho^{(\text{ind})}(\vec{r})$, the electric field $\vec{E}(\vec{r})$, and the displacement field $\vec{D}(\vec{r})$ for the static case. Introduce an appropriate limit to recover the properties of an ideal conductor.

Hint: Since the question involves a calculation from Problem 4.3, you may refer to this result.

c) Neglecting dissipation, $\tau \to \infty$, the material supports longitudinal and transverse waves. Discuss the monochromatic plane waves; in particular, derive the dispersion relation of the longitudinal plasma oscillations and the transverse modes.