

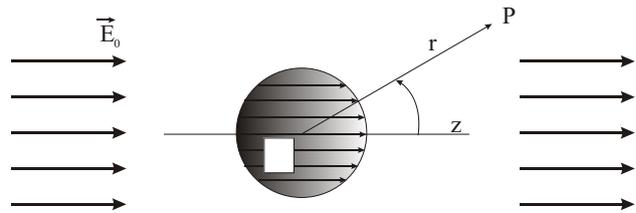


**T II: Elektrodynamik**  
(Prof. E. Frey)

**Problem set 10**

**Tutorial 10.1** Dielectric sphere

A dielectric sphere of radius  $R$  characterized by a dielectric constant  $\epsilon$  is placed in an initially uniform electric field  $\vec{E}_0$ . For convenience, choose the center of the sphere as the origin and consider  $\vec{E}_0$  along the  $z$ -axis. Then the problem exhibits an axial symmetry, which simplifies the problem.



- a) Determine the electrostatic potential  $\varphi(\vec{x})$  inside the sphere,  $|\vec{x}| \leq R$ , and outside the sphere,  $|\vec{x}| > R$ . Formulate appropriate matching conditions at the surface of the sphere. Recall that the most general axially symmetric solution of Laplace's equation  $\nabla^2 \varphi = 0$  in polar coordinates,  $\vec{x} = r(\sin \vartheta \cos \phi, \sin \vartheta \sin \phi, \cos \vartheta)$ , is given in terms of

$$\varphi(\vec{x}) = \sum_{\ell=0}^{\infty} \left( a_{\ell} r^{\ell} + b_{\ell} r^{-(\ell+1)} \right) P_{\ell}(\cos \vartheta),$$

where  $P_{\ell}(t)$  denotes the Legendre polynomials and  $a_{\ell}, b_{\ell}$  are undetermined coefficients.

- b) Derive the corresponding electric field  $\vec{E}(\vec{x})$ . Extract the polarization  $\vec{P}(\vec{x})$  inside the sphere and find the total induced dipole moment  $\vec{p}$  of the dielectric sphere. Determine the effective polarizability of the sphere  $\alpha$  defined by  $\vec{p} = \alpha \vec{E}_0$ .
- c) Show that charges accumulate at the surface of the sphere and determine the induced surface charge density  $\sigma$ .

**Problem 10.2**     *Magnetic shielding –  $\mu$ -metal*

A  $\mu$ -metal is a nickel-iron alloy that has a very high magnetic permeability  $\mu \sim 10^4 - 10^6 \gg 1$ . The technical application of these materials is the screening of static (or low-frequency) magnetic fields, which cannot be attenuated by other methods.

Consider a spherical shell of magnetic permeability  $\mu$  and inner and outer radii  $R_i$  and  $R_o$ , respectively, placed in a previously uniform magnetic field  $\vec{H}_\infty$ . The medium inside and outside of the shell has a magnetic permeability  $\mu = 1$ .

- Argue that one can introduce a scalar magnetic potential  $\varphi_M$  to represent the field  $\vec{H} = -\vec{\nabla}\varphi_M$ , and show that it fulfills the Laplace equation  $\nabla^2\varphi_M = 0$  in each region.
- State the appropriate matching conditions for  $\varphi_M$  at the interfaces  $r = R_i$  and  $r = R_o$ .
- Recalling that the most general solution of the Laplace equation with cylindrical symmetry is provided by

$$\varphi_M(\vec{r}) = \sum_{\ell=0}^{\infty} \left( a_\ell r^\ell + b_\ell r^{-(\ell+1)} \right) P_\ell(\cos \vartheta),$$

determine the magnetostatic potential in each region. Calculate explicitly the corresponding magnetic field  $\vec{H}$  inside of the shell.

- Determine the leading behavior of the field inside of a thin shell of a  $\mu$ -metal for  $\mu \rightarrow \infty$ . Discuss why a  $\mu$ -metal provides an effective shielding.

**Problem 10.3**     *Legendre polynomials*

Consider the following partial differential equation

$$\frac{\partial}{\partial t} \left[ (1-t^2) \frac{\partial \psi}{\partial t} \right] = -\frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right). \quad (*)$$

A solution is provided by the function

$$\psi(r, t) = \frac{1}{\sqrt{1-2rt+r^2}}, \quad -1 < r < 1, \quad -1 \leq t \leq 1,$$

which serves as a generating function for the *Legendre polynomials*  $P_\ell(t)$ , i.e., a Taylor expansion with respect to  $r$ ,  $\psi(r, t) = \sum_{\ell=0}^{\infty} r^\ell P_\ell(t)$ , defines the functions  $P_\ell(t)$ .

- Show by explicit substitution that  $\psi(r, t)$  indeed solves the partial differential equation (\*).
- Identifying  $t = \cos \vartheta$  reveals that  $\psi(r, t)$  corresponds to the Coulomb potential of a unit charge located on the  $z$ -axis at unit distance from the origin. Thus  $\psi(r, t = \cos \vartheta)$  solves the Laplace equation in polar coordinates

$$\nabla^2 \psi(r, \cos \vartheta) = \frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \vartheta} \frac{\partial}{\partial \vartheta} \left( \sin \vartheta \frac{\partial \psi}{\partial \vartheta} \right) = 0.$$

Using this observation derive the partial differential equation (\*).

- Determine explicitly  $P_\ell(t = \pm 1)$  observing that for  $t = \pm 1$  the Taylor series of  $\psi(r, t)$  becomes elementary.
- Employ the symmetries of  $\psi(r, t)$  to argue that  $P_\ell(t)$  is a symmetric (anti-symmetric) function for even (odd)  $\ell$ .
- Inspect the Taylor series to demonstrate that  $P_\ell(t)$  is a polynomial of order  $\ell$ .  
*Hint:* Expand the square root in  $x = 2rt - r^2$ .
- Calculate and sketch the first four Legendre polynomials ( $\ell = 0, \dots, 3$ ).

- g) Substitute the Taylor series of  $\psi(r, t)$  in the partial differential equation (\*). Comparing the coefficients of  $r^\ell$  confirm that the  $P_\ell(t)$  satisfy the second order differential equation, i.e., they are indeed Legendre polynomials,

$$\frac{d}{dt} \left[ (1-t^2) \frac{dP_\ell(t)}{dt} \right] + \ell(\ell+1)P_\ell(t) = 0, \quad -1 \leq t \leq 1. \quad (**)$$

- h) Show that  $(1-2rt+r^2)\partial\psi/\partial r = (t-r)\psi$ . Make use of this result to derive the recursion relation

$$\ell P_{\ell-1}(t) - (2\ell+1)tP_\ell(t) + (\ell+1)P_{\ell+1}(t) = 0.$$

Similarly, verify that  $(1-2rt+r^2)\partial\psi/\partial t = r\psi$  and prove

$$P'_{\ell+1}(t) - 2tP'_\ell(t) + P'_{\ell-1}(t) = P_\ell(t).$$

- i\*) Show that the Legendre polynomials are orthogonal in the following sense,

$$\int_{-1}^1 dt P_\ell(t)P_{\ell'}(t) = \frac{2}{2\ell+1}\delta_{\ell\ell'}.$$

*Hint:* Employ the differential equation (\*\*) to show that

$$[\ell'(\ell'+1) - \ell(\ell+1)] \int_{-1}^1 dt P_\ell(t)P_{\ell'}(t) = 0$$

and conclude that orthogonality holds. The normalization follows by considering  $\int dt \psi(r, t)^2$ . First perform the integration directly; then use the Taylor expansion in  $r$  and the orthogonality property.

*Due date: Tuesday, 7/3/07, at 9 p.m.*