

T IV: Thermodynamik und Statistik
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Problem set 5

Problem 5.1 *density matrix*

For many quantum problems it is sufficient to consider only two states, by analogy to a spin 1/2 referred to as 'spin up' and 'spin down'. The corresponding Hilbert space \mathcal{E} is then two-dimensional and operators acting on \mathcal{E} can be represented by 2×2 matrices. Show that any operator is represented by a linear combination of the identity matrix \mathbb{I} and the Pauli matrices

$$\sigma_1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad \sigma_2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad \sigma_3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$

Construct the normed eigenstates $|\sigma k\rangle$, $k = 1, 2, 3$ of the Pauli matrices and demonstrate that the eigenvalues σ are given by ± 1 . Construct the projectors $\Lambda_{k+} = |k\rangle\langle k|$ and consider density matrices $\rho = \sum_{k=1}^3 q_k \Lambda_{k+}$ with $\sum_k q_k = 1$.

1. Express the projectors Λ_{k+} in terms of the identity and the Pauli matrices. Show that the density matrix can be written as

$$\rho = \frac{1}{2} (\mathbb{I} + \mathbf{P} \cdot \boldsymbol{\sigma}), \quad \mathbf{P} = \langle \boldsymbol{\sigma} \rangle = \text{Tr}(\boldsymbol{\sigma} \rho)$$

2. For the density matrix ρ a measurement of the observable σ_1 is performed. Calculate the probability W_{1+} that the system is found in eigenstate $|+1\rangle$ after the measurement. Interpret the result in terms of the probabilities q_k .
3. Determine the expectation and variance for measuring $\mathbf{n} \cdot \boldsymbol{\sigma}$, i.e. $\langle \mathbf{n} \cdot \boldsymbol{\sigma} \rangle$ and $\langle (\mathbf{n} \cdot \boldsymbol{\sigma})^2 \rangle - \langle \mathbf{n} \cdot \boldsymbol{\sigma} \rangle^2$ for arbitrary unit vectors \mathbf{n} .
4. Evaluate the probabilities $W_{\mathbf{n}\pm}$ to find eigenvalues ± 1 when measuring the observable $\mathbf{n} \cdot \boldsymbol{\sigma}$ for the density matrix ρ .
5. Determine the density matrix that maximizes the functional $S[\rho] = -\text{Tr}(\rho \ln \rho)$ with the constraint that probability is preserved $\text{Tr} \rho = 1$ and the average $\mathbf{P} = \langle \boldsymbol{\sigma} \rangle$ is known. In general \mathbf{P} can have a norm smaller than unity. Use appropriate Lagrange multipliers to enforce the constraints.

Problem 5.2 *spin precession*

A single spin described by density matrix ρ is exposed to a constant magnetic field. The Hamiltonian for the problem is given by $\mathcal{H} = -(\gamma\hbar/2)\mathbf{B} \cdot \boldsymbol{\sigma}$, where $\boldsymbol{\sigma} = (\sigma_1, \sigma_2, \sigma_3)$ are the Pauli matrices. Use the von Neumann equations to show that the magnetization $\mathbf{P}(t) = (\gamma\hbar/2)\text{Tr}\rho(t)\boldsymbol{\sigma}$ obeys the classical equation of motion

$$\dot{\mathbf{P}}(t) = -\mathbf{B} \times \mathbf{P}(t)$$

* Assume a collection of spins in an external magnetic field \mathbf{B}_0 . Additionally there may be small random local fields, such that the Larmor frequency $\omega = \gamma B$ is statistically distributed. Assuming a Gaussian or Lorentzian distribution for ω characterized by mean $\omega_0 = \gamma B_0$ and variance $1/T^2$ derive the time dependence of the average magnetization $\overline{\mathbf{P}}(t)$ by averaging the solution of the classical equation of motion for a single spin and show that it quickly loses phase coherence.

Problem 5.3 *entangled states*

The density matrix of a subsystem is usually in a mixed state even if the entire system is in a pure state. Assuming that the wave function of the entire system is merely a product of wave functions for the subsystem and the reservoir, show that the density matrix of the subsystem corresponds to a pure state.

If the total system is described by a pure state and the density matrix of the subsystem describes a pure state, show that the wave function of the total system necessarily factorizes.

Can you construct a mixed state for the density matrix such that the subsystem is in a pure state?

Problem 5.4 *Pressure ensemble*

Use the NPT (constant pressure) ensemble with the phase space density

$$\rho_P = Z_P^{-1} \exp(-\beta\mathcal{H} - \beta PV), \quad Z_P(T, P, N) = \int_0^\infty dV \int \frac{d^{3N}r d^{3N}p}{N!h^{3N}} \exp(-\beta\mathcal{H} - \beta PV)$$

and the definition of averages to show that

$$Z_P(T, P, N) = Z_P(T, P_0, N) \langle e^{\beta(P_0 - P)V} \rangle_0.$$

Here $\langle \cdot \rangle_0$ denotes averaging with respect to ρ_P at pressure P_0 . Derive the corresponding expansion of the free enthalpy $G(T, P, N) = -k_B T \ln Z_P(T, P, N)$ in terms of the cumulants of the volume.

In the thermodynamic limit, i.e. $N \rightarrow \infty$, P and T fixed, derive the thermodynamic relations

$$G(T, P, N) = F(T, \langle V \rangle, N) + P\langle V \rangle, \quad dG = -SdT + \langle V \rangle dP + \mu dN$$