

David
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Introduction

β -supergravity

BI, NS-branes, ...

Conclusion

Vacua of β -supergravity

David ANDRIOT

AEI Potsdam-Golm, Germany

arXiv:1306.4381, 1402.5972 and work in progress

by D. A. and André Betz

Frontiers in String Phenomenology,
01/08/2014, Ringberg Castle, Tegernsee, Germany

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NSNS sector in 10D and 4D supergravities

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Introduction

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4D gauged supergravity: also non-geometric fluxes $Q_a{}^{bc}, R^{abc}$

all fluxes are components of the embedding tensor

\Leftrightarrow some “structure constants” in gauging algebra

hep-th/0508133 by J. Shelton, W. Taylor, B. Wecht

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\Rightarrow specific terms in 4D super/scalar potential, \checkmark for pheno !

Moduli stabilisation, de Sitter solutions with $Q_a{}^{bc}, R^{abc}$:

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Terms/fluxes related by 4D T-duality chain

$$H_{abc} \xleftrightarrow{T_a} f^a{}_{bc} \xleftrightarrow{T_b} Q_c{}^{ab} \xleftrightarrow{T_c} R^{abc}$$

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10D origin of Q and R ? Naively not present in $\mathcal{L}_{\text{NSNS}}$

+ argued to descend from 10D non-geometric backgrounds...

Original idea of 10D non-geometry

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A (target) space, divided in patches

Fields glue on overlaps:

diffeomorphisms, gauge transformation

(point-like symmetries)

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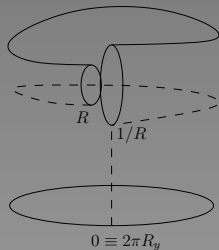
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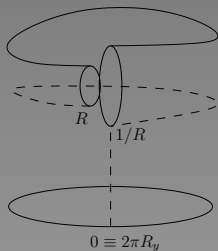
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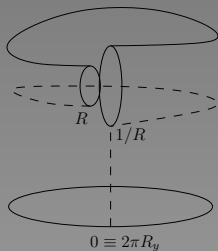
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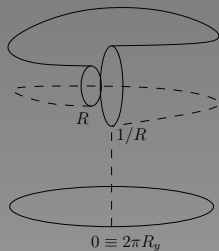
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β -supergravity: reformulation of $\mathcal{L}_{\text{NSNS}}$ after field redefinition

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Conclusion

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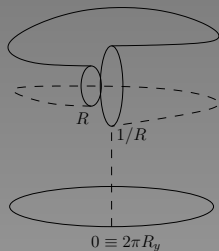
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↪ $\tilde{\mathcal{L}}_{\beta}(f^a{}_{bc}, Q_a{}^{bc}, R^{abc})$, and standard geometry restored

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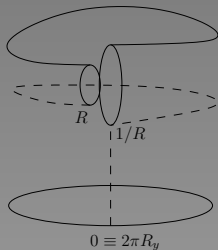
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↪ $\tilde{\mathcal{L}}_{\beta}(f^a{}_{bc}, Q_a{}^{bc}, R^{abc})$, and standard geometry restored

↪ allows compactification, uplift of 4D gauged supergravities.

10D Bianchi identities for fluxes:

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Introduction

β -supergravity

BI, *NS*-branes, ...

Conclusion

10D Bianchi identities for fluxes:

$$dH = 0 \Rightarrow \partial_{[a} H_{bcd]} - \frac{3}{2} f^e_{[ab} H_{cd]e} = 0$$

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Correction in presence of source: *NS5*-brane: $dH \sim \text{vol}_4 \delta^{(4)}(r_4)$
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\mathcal{D} is the Dirac operator associated to the $Spin(d, d) \times \mathbb{R}^+$
covariant derivative (Generalized Geometry, DFT).

[arXiv:1107.0008](#) by O. Hohm, S. K. Kwak and B. Zwiebach

[arXiv:1304.1472](#) by D. Geissbühler, D. Marqués, C. Núñez, V. Penas

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Show completely analogous results in β -supergravity for Q, R ;
 Q -flux is sourced by 5_2^2 - or Q -brane.

\mathcal{D} is the Dirac operator associated to the $Spin(d, d) \times \mathbb{R}^+$
covariant derivative (Generalized Geometry, DFT).

[arXiv:1107.0008](#) by O. Hohm, S. K. Kwak and B. Zwiebach

[arXiv:1304.1472](#) by D. Geissbühler, D. Marqués, C. Núñez, V. Penas

Relation to conditions for supersymmetric vacua with
 $SU(3) \times SU(3)$ structure.

β -supergravity

Lagrangian $\tilde{\mathcal{L}}_\beta(f, Q, R)$

Reformulation of $\mathcal{L}_{\text{NSNS}}$ via a field redefinition

Build on earlier results:

[arXiv:1303.0251](#) by D. A.

[arXiv:1106.4015](#) by D. A., M. Larfors, D. Lüst, P. Patalong

[arXiv:1202.3060](#), [arXiv:1204.1979](#) by D. A., O. Hohm, M. Larfors, D. Lüst, P. Patalong

See also related work in:

[arXiv:1210.1591](#), [arXiv:1211.0030](#), [arXiv:1304.2784](#)

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Idea: field redef. $(g_{mn}, b_{mn}, \phi) \leftrightarrow (\tilde{g}_{mn}, \beta^{mn}, \tilde{\phi})$, β antisym.

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$$\begin{aligned} \tilde{g}^{-1} &= (g + b)^{-1} g (g - b)^{-1} \\ \beta &= -(g + b)^{-1} b (g - b)^{-1} \end{aligned} \Leftrightarrow (g + b)^{-1} = (\tilde{g}^{-1} + \beta), \quad \frac{e^{-2\tilde{\phi}}}{e^{-2\phi}} = \frac{\sqrt{|g|}}{\sqrt{|\tilde{g}|}}$$

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$$\mathcal{H} = \begin{pmatrix} g - bg^{-1}b & -bg^{-1} \\ g^{-1}b & g^{-1} \end{pmatrix} = \mathcal{E}^T \mathbb{I} \mathcal{E} = \tilde{\mathcal{E}}^T \mathbb{I} \tilde{\mathcal{E}} = \begin{pmatrix} \tilde{g} & \tilde{g}\beta \\ -\beta\tilde{g} & \tilde{g}^{-1} - \beta\tilde{g}\beta \end{pmatrix}$$

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β -supergravity

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$\tilde{g} = \tilde{e}^T \eta_d \tilde{e}$
 β w.r.t. Q, R : motivations from Gen. Complex Geom./sugra

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Apply the field redefinition to $\mathcal{L}_{\text{NSNS}}(g, b, \phi)$

Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Apply the field redefinition to $\mathcal{L}_{\text{NSNS}}(g, b, \phi) = \tilde{\mathcal{L}}_{\beta}(\tilde{g}, \beta, \tilde{\phi}) + \partial(\dots)$

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To see the fluxes, Q , better to use flat tangent space indices

$$Q_c{}^{ab} = \partial_c \beta^{ab} - 2\beta^{d[a} f^b]{}_{cd}, \quad R^{abc} = 3\beta^{d[a} \nabla_d \beta^{bc]}$$

arXiv:0807.4527 by M. Graña, R. Minasian, M. Petrini, D. Waldram

arXiv:1109.0290 by G. Aldazabal, W. Baron, D. Marqués, C. Núñez

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Nice structure w.r.t. 4D, with $Q_a{}^{ab} = 0$

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10D theory with non-geometric fluxes (β -supergravity),
uplift of 4D gauged supergravity ✓

More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, *NS*-branes, ...

Conclusion

More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

- For standard ∇_m and Levi-Civita connection Γ_{np}^m

$$\tilde{e}^a{}_m \tilde{e}^n{}_b \nabla_n V^m = \nabla_b V^a \equiv \partial_b V^a + \omega_{bc}^a V^c$$

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More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

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$$\check{\nabla}^m V^p = -\beta^{mn} \partial_n V^p - \check{\Gamma}_n^{mp} V^n, \quad \check{\nabla}^m V_p = -\beta^{mn} \partial_n V_p + \check{\Gamma}_p^{mn} V_n$$

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More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

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More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

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More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

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Define as well a new "Ricci tensor and scalar"

\hookrightarrow enters in Lagrangian in curved/flat indices.

More structure

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

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$$\Leftrightarrow -\omega_{Q_a}{}^{bc} \equiv \tilde{e}^b{}_n \tilde{e}^m{}_a \left(-\beta^{nq} \partial_q \tilde{e}^c{}_m + \tilde{e}^c{}_p \check{\Gamma}_m^{np} \right)$$

$$\omega_{Q_a}{}^{bc} = \frac{1}{2} \left(Q_a{}^{bc} + \eta_{ad} \eta^{ce} Q_e{}^{db} + \eta_{ad} \eta^{be} Q_e{}^{dc} \right) !!$$

Define as well a new "Ricci tensor and scalar"

\hookrightarrow enters in Lagrangian in curved/flat indices.

Structures: natural with Generalized Geometry formalism

10D examples of non-geometries

David
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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

10D examples of non-geometries

David
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Famous toroidal example: NSNS sector:

$$T^3 + H_{123} \xleftrightarrow{T_1} \text{twisted torus } (f^1_{23}) \xleftrightarrow{T_2} \text{non-geometric config.}$$

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

10D examples of non-geometries

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

10D examples of non-geometries

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

10D examples of non-geometries

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

10D examples of non-geometries

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

10D examples of non-geometries

David
ANDRIOT

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

David
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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

David
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Field redefinition always possible off-shell/locally

$$\mathcal{L}_{\text{NSNS}}(g, b, \phi) = \tilde{\mathcal{L}}_{\beta}(\tilde{g}, \beta, \tilde{\phi}) + \partial(\dots)$$

Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

David
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Geometry restored: T^3 , Q -flux \checkmark .

Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

David
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Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

David
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True for a class of non-geometric/geometric backgrounds:

n isom., glue with " β -transforms" $\in O(n, n)$ and diffeos.

arXiv:1402.5972 by D. A. and André Betz

Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

Non-geometry/geometry: global aspects

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β -supergravity has 10D non-geometric fluxes,
it restores geometry \Rightarrow alternative description, compactif. \checkmark

Introduction

β -supergravity

Lagrangian

More structure

(Non)-geometry

BI, NS-branes, ...

Conclusion

David
ANDRIOT

Bianchi identities and NS -branes

Bianchi identities for the fluxes

Introduction

β -supergravity

BI, NS -branes, ...

BI

NS -branes

\mathcal{D}

SUSY

Conclusion

Bianchi identities and *NS*-branes

Bianchi identities for the fluxes

10D standard supergravity: $2\partial_{[a}H_{bcd]} - 3f^e{}_{[ab}H_{cd]e} = 0$

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Introduction

β -supergravity

BI, *NS*-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

Bianchi identities and NS-branes

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4D, constant fluxes: gaugings of gauged supergravity

$$[Z_a, Z_b] = H_{abc}X^c + f^c{}_{ab}Z_c$$

$$[Z_a, X^b] = -f^b{}_{ac}X^c + Q_a{}^{bc}Z_c$$

$$[X^a, X^b] = Q_c{}^{ab}X^c - R^{abc}Z_c$$

[hep-th/0508133](#) by J. Shelton, W. Taylor, B. Wecht

Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

Bianchi identities and NS-branes

Bianchi identities for the fluxes

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Jacobi identities \Leftrightarrow Bianchi identities

$$3f^e{}_{[ab}H_{cd]e} = 0$$

$$H_{d[ab}Q_f]{}^{ed} + f^e{}_d{}[af^d{}_{bf}] = 0$$

$$\frac{1}{2}H_{gaf}R^{deg} - \frac{1}{2}Q_g{}^{de}f^g{}_{af} + 2Q_{[a}{}^{g[d}f^e]{}_{f]g} = 0$$

$$3R^{d[gh}f^i]{}_{ad} - 3Q_a{}^{d[gh}Q_d{}^{hi]} = 0$$

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10D realisation in β -supergravity (non-constant fluxes), $H = 0$

Bianchi identities and NS-branes

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10D standard supergravity: $2\partial_{[a}H_{bcd]} - 3f^e{}_{[ab}H_{cd]e} = 0$
 $\partial_{[b}f^a{}_{cd]} - f^a{}_e[bf^e{}_{cd]} = 0$

4D, constant fluxes: gaugings of gauged supergravity

$$[Z_a, Z_b] = H_{abc}X^c + f^c{}_{ab}Z_c$$

$$[Z_a, X^b] = -f^b{}_{ac}X^c + Q_a{}^{bc}Z_c$$

$$[X^a, X^b] = Q_c{}^{ab}X^c - R^{abc}Z_c$$

[hep-th/0508133](https://arxiv.org/abs/hep-th/0508133) by J. Shelton, W. Taylor, B. Wecht

Jacobi identities \Leftrightarrow Bianchi identities

$$\partial_{[a}f^e{}_{bf]} - f^e{}_d[a f^d{}_{bf]} = 0$$

$$\partial_{[a}Q_{f]}{}^{de} - \beta^{g[d}\partial_g f^e]{}_{af} - \frac{1}{2}Q_g{}^{de}f^g{}_{af} + 2Q_{[a}{}^g[d f^e]{}_{f]g} = 0$$

$$\partial_a R^{ghi} - 3\beta^{d[g}\partial_d Q_a{}^{hi]} + 3R^{d[gh}f^i]{}_{ad} - 3Q_a{}^{d[g}Q_d{}^{hi]} = 0$$

$$2\beta^{g[d}\partial_g R^{abc]} + 3R^{g[da}Q_g{}^{bc]} = 0$$

10D realisation in β -supergravity (non-constant fluxes), $H = 0$

Bianchi identities and NS-branes

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10D realisation in β -supergravity (non-constant fluxes), $H = 0$
 Automatically satisfied with 10D expressions of f, Q, R

Bianchi identities and NS-branes

Bianchi identities for the fluxes

10D standard supergravity: $2\partial_{[a}H_{bcd]} - 3f^e{}_{[ab}H_{cd]e} = 0$
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Jacobi identities \Leftrightarrow Bianchi identities

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10D realisation in β -supergravity (non-constant fluxes), $H = 0$
 Automatically satisfied with 10D expressions of f, Q, R
 More involved expressions for Bianchi identities

NS-branes as sources

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Introduction

β -supergravity

BI, *NS*-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

NS-branes as sources

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Introduction

β -supergravity

BI, *NS*-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

$$NS5\text{-brane} \begin{array}{c} \xleftarrow{\text{smearing}} \\ \xrightarrow{+T\text{-d.}} \end{array} KK\text{-monopole} \begin{array}{c} \xleftarrow{\text{smearing}} \\ \xrightarrow{+T\text{-d.}} \end{array} 5_2^2 \text{ or } Q\text{-brane}$$

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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$$NS5\text{-brane: } dH \sim \text{vol}_4 \delta^{(4)}(r_4)$$

$$\partial_{[a} H_{bcd]} - \frac{3}{2} f^e{}_{[ab} H_{cd]e} = \frac{c_H}{4} \epsilon_{4\perp abcd} \delta^{(4)}(r_4)$$

$$\hookrightarrow \text{Poisson equation: } \Delta_4 f_H = c_H \delta^{(4)}(r_4)$$

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Source corrections to Bianchi identities

NS-branes as sources

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Source corrections to Bianchi identities

KK-monopole:

$$\partial_{[b} f^a{}_{cd]} - f^a{}_{e[b} f^e{}_{cd]} = \frac{C_K}{3} \epsilon_{3\perp bcd} \epsilon_{1||e} \eta^{ea} \delta^{(3)}(r_3)$$

see also

[arXiv:0706.3049](https://arxiv.org/abs/0706.3049) by G. Villadoro and F. Zwirner

NS-branes as sources

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Q-brane:

$$\begin{aligned} \partial_{[a} Q_{b]}{}^{cd} - \beta g^{[c} \partial_g f^{d]}{}_{ab} - \frac{1}{2} Q_g{}^{cd} f^g{}_{ab} + 2 Q_{[a}{}^{g[c} b^{d]}{}_{f]g} \\ = \frac{C_Q}{2} \epsilon_{2\perp ab} \epsilon_{2||ef} \eta^{ec} \eta^{fd} \delta^{(2)}(r_2) \end{aligned}$$

NS-branes as sources

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Source corrections to Bianchi identities

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On brane solutions \Rightarrow Poisson equations

$$KK\text{-monopole: } \Delta_3 f_K = c_K \delta^{(3)}(r_3)$$

$$Q\text{-brane: } \Delta_2 f_Q = c_Q \delta^{(2)}(r_2)$$

Generalized derivative \mathcal{D}

10D standard supergravity: $\mathcal{D}^2 = 0 \Leftrightarrow$ Bianchi identities

$$\mathcal{D}A = 2e^\phi(d - H \wedge)(e^{-\phi}A)$$

Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

Generalized derivative \mathcal{D}

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10D standard supergravity: $\mathcal{D}^2 = 0 \Leftrightarrow$ Bianchi identities

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$$= 2 \left(\partial_a \cdot e^a \wedge - \frac{1}{2} f^c{}_{ab} e^a \wedge e^b \wedge \iota_c - \frac{1}{6} H_{abc} e^a \wedge e^b \wedge e^c \wedge - \partial_a \phi e^a \wedge \right) A$$

$$\text{where } V \lrcorner A = V^a \iota_a A = \frac{1}{(p-1)!} V^{m_1} A_{m_1 \dots m_p} dx^{m_2} \wedge \dots \wedge dx^{m_p}$$

Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

Generalized derivative \mathcal{D}

David
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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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In 4D:

hep-th/0607015 by J. Shelton, W. Taylor, B. Wecht

arXiv:0705.3410 by M. Ihl, D. Robbins, and T. Wrase

$$\begin{aligned} \mathcal{D}_{\sharp} A = & \left(-\frac{1}{6} H_{abc} \tilde{e}^a \wedge \tilde{e}^b \wedge \tilde{e}^c \wedge - \frac{1}{2} f^a{}_{bc} \tilde{e}^b \wedge \tilde{e}^c \wedge \iota_a \right. \\ & - \frac{1}{2} Q_c{}^{ab} \tilde{e}^c \wedge \iota_a \iota_b + \frac{1}{6} R^{abc} \iota_a \iota_b \iota_c \\ & \left. - \frac{1}{2} f^a{}_{ab} \tilde{e}^b \wedge + \frac{1}{2} Q_a{}^{ab} \iota_b \right) A \end{aligned}$$

Generalized derivative \mathcal{D}

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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$$\mathcal{D}_{\sharp}^2 = 0 \Leftrightarrow \text{Bianchi identities and } \frac{1}{3}H_{abc}R^{abc} + \frac{1}{2}f^a{}_{ab}Q_a{}^{ab} = 0$$

Generalized derivative \mathcal{D}

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Realisation in 10D, for non-geometric fluxes, in β -supergravity:

hep-th/0607015 by J. Shelton, W. Taylor, B. Wecht
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Generalized derivative \mathcal{D}

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Realisation in 10D, for non-geometric fluxes, in β -supergravity:

$$\mathcal{D} = 2\mathcal{D}_{\sharp} + 2 \text{ other pieces}$$

Generalized derivative \mathcal{D}

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

10D standard supergravity: $\mathcal{D}^2 = 0 \Leftrightarrow$ Bianchi identities

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Realisation in 10D, for non-geometric fluxes, in β -supergravity:

$\mathcal{D} = 2\mathcal{D}_\sharp + 2$ other pieces

$$\mathcal{D} = 2\left(\partial_a \cdot \tilde{e}^a \wedge + \beta^{ab}\partial_b \cdot \iota_a - \frac{1}{2}f^c{}_{ab}\tilde{e}^a \wedge \tilde{e}^b \wedge \iota_c - \frac{1}{2}Q_a{}^{bc}\tilde{e}^a \wedge \iota_b \iota_c \right. \\ \left. + Q_d{}^{dc}\iota_c + \frac{1}{6}R^{abc}\iota_a \iota_b \iota_c - \partial_a \tilde{\phi} \tilde{e}^a \wedge - (\beta^{ab}\partial_b \tilde{\phi} - \mathcal{T}^a)\iota_a \right)$$

Generalized derivative \mathcal{D}

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

10D standard supergravity: $\mathcal{D}^2 = 0 \Leftrightarrow$ Bianchi identities

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$$= 2\left(\partial_a \cdot e^a \wedge - \frac{1}{2}f^c{}_{ab}e^a \wedge e^b \wedge \iota_c - \frac{1}{6}H_{abc}e^a \wedge e^b \wedge e^c \wedge - \partial_a \phi e^a \wedge\right)A$$

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Realisation in 10D, for non-geometric fluxes, in β -supergravity:

$\mathcal{D} = 2\mathcal{D}_{\sharp} + 2$ other pieces

$$\begin{aligned} \mathcal{D} = & 2\left(\partial_a \cdot \tilde{e}^a \wedge + \beta^{ab}\partial_b \cdot \iota_a - \frac{1}{2}f^c{}_{ab}\tilde{e}^a \wedge \tilde{e}^b \wedge \iota_c - \frac{1}{2}Q_a{}^{bc}\tilde{e}^a \wedge \iota_b \iota_c \right. \\ & \left. + Q_d{}^{dc}\iota_c + \frac{1}{6}R^{abc}\iota_a \iota_b \iota_c - \partial_a \tilde{\phi} \tilde{e}^a \wedge - (\beta^{ab}\partial_b \tilde{\phi} - \mathcal{T}^a)\iota_a\right) \end{aligned}$$

$$\mathcal{D}A = 2e^{\tilde{\phi}}(\nabla_a \cdot \tilde{e}^a \wedge - \tilde{\nabla}^a \cdot \iota_a + \mathcal{T} \vee + R \vee)(e^{-\tilde{\phi}}A)$$

Generalized derivative \mathcal{D}

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Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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In 4D:

hep-th/0607015 by J. Shelton, W. Taylor, B. Wecht

arXiv:0705.3410 by M. Ihl, D. Robbins, and T. Wrase

$$\begin{aligned} \mathcal{D}_{\sharp}A = & \left(-\frac{1}{6}H_{abc}\tilde{e}^a \wedge \tilde{e}^b \wedge \tilde{e}^c \wedge - \frac{1}{2}f^a{}_{bc}\tilde{e}^b \wedge \tilde{e}^c \wedge \iota_a \right. \\ & - \frac{1}{2}Q_c{}^{ab}\tilde{e}^c \wedge \iota_a \iota_b + \frac{1}{6}R^{abc}\iota_a \iota_b \iota_c \\ & \left. - \frac{1}{2}f^a{}_{ab}\tilde{e}^b \wedge + \frac{1}{2}Q_a{}^{ab}\iota_b \right)A \end{aligned}$$

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David
ANDRIOT

Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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David
ANDRIOT

Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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Γ^A : Clifford algebra: $\{\Gamma^A, \Gamma^B\} = 2\eta^{AB}$, use the representation

$\Gamma^A : \Gamma^a = 2\tilde{e}^a \wedge, \Gamma_a = 2\iota_a$, spinor Ψ : polyform

SUSY vacua with $SU(3) \times SU(3)$ structure

For standard supergravity, $\mathcal{D}A = 2e^\phi(d - H \wedge)(e^{-\phi}A)$ appears for RR (spinors): $F \sim \mathcal{D}C$

Introduction

β -supergravity

BI, NS-branes, ...

BI

NS-branes

\mathcal{D}

SUSY

Conclusion

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where $e^A = \text{constant}$, $|\mu|^2 \sim -\Lambda$, and

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[hep-th/0505212](#), [hep-th/0609124](#) by M. Graña, R. Minasian, M. Petrini and A. Tomasiello

For an $SU(3)$ structure: $\Phi_+ \sim e^{-iJ}$, $\Phi_- \sim \Omega_3$.

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Study superpotential, and further properties of $\mathcal{D}...$

Conclusion

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4D non-geometric fluxes, 10D non-geometry

Introduction

β -supergravity

BI, NS-branes, ...

Conclusion

Conclusion

David
ANDRIOT

Introduction

β -supergravity

BI, NS-branes, ...

Conclusion

4D non-geometric fluxes, 10D non-geometry

β -supergravity: 10D theory with non-geometric fluxes Q, R
non-geometric bckgd of standard supergravity \rightarrow geometric
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Conclusion

David
ANDRIOT

Introduction

β -supergravity

BI, NS -branes, ...

Conclusion

4D non-geometric fluxes, 10D non-geometry

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10D Bianchi identities, relation to 4D Jacobi identities

Source corrections with NS -branes, Poisson equations.

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Conclusion

David
ANDRIOT

Introduction

β -supergravity

BI, *NS*-branes, ...

Conclusion

4D non-geometric fluxes, 10D non-geometry

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Generalized Geometry/Double Field Theory: formalisms that
reproduce standard and β -supergravity; structures: natural.

$Spin(d, d) \times \mathbb{R}^+$ cov. derivative, Lagrangian, e.o.m., SUSY.

Conclusion

David
ANDRIOT

Introduction

β -supergravity

BI, NS-branes, ...

Conclusion

4D non-geometric fluxes, 10D non-geometry

β -supergravity: 10D theory with non-geometric fluxes Q, R
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Extension beyond the NSNS sector: RR non-geometric fluxes?
exotic D -branes? Exceptional geometry/field theory could help
Additional Bianchi identities $\Rightarrow R$ -brane?

\leftrightarrow Get new pheno. interesting backgrounds, de Sitter vacua.