

Asymmetric CFTs and GSUGRA II

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based on 1608.00595 and 1611.04617 by R. Blumenhagen, MF, E. Plauschinn

November 24, 2016

General Idea

Two separate ways to stabilize moduli:

- **SUGRA (TS): Fluxes/gaugings**
- **CFT (WS): (Asymmetric) Orbifolds**

Work on L-R asymmetric torodial orbifolds suggests a connection:
After introducing the asymmetry one finds a flux algebra!

[Dabholkar, Hull '02,05; Condeescu, Florakis, Kounnas, Lüst '12,13]

GSUGRA \sim ACFT?!

Overview

What we did:

Look at Gepner models + L-R asymmetric simple currents

[Gepner; Schellekens, Yankielowicz; Schellekens, Gato-Rivera]

Compare the result to a SUGRA with NSNS gaugings

Two papers together with R. Blumenhagen and E. Plauschinn:

- 1608.00595 Very concrete examples in 4D with $\mathcal{N} = 1$ SUSY
- 1611.04617 Classification of asymmetric Gepner models in 4D, 6D, 8D with extended SUSY to support conjecture.

Our results suggest: Yes! GSUGRA \sim ACFT !

Recap: The 3^5 Gepner model

Gepners idea: Use tensored minimal SCFTs as the internal CFT of a string compactification.

Example: Take the CFT $(k=3)^5$ to describe a 6D internal space. The massless states look like e.g.

$$(\mathbf{3}, 4, 1)(\mathbf{2}, 3, 1)(\mathbf{0}, 1, 1)^3 C \rightarrow x_1^3 x_2^2$$

$$(\mathbf{2}, 3, 1)(\mathbf{1}, 2, 1)^3(\mathbf{0}, 1, 1) C \rightarrow x_1^2 x_2 x_3 x_4$$

- ⇒ Combinatorics of complex structure deformations in $\mathbb{P}_{1,1,1,1,1}[5]$.
- ⇒ 3^5 model is IIB on the quintic at a certain point in moduli space.
- ⇒ $\mathcal{N} = 2$ target space SUSY.

In general: More complicated WCP

Now: Add a certain L-R asymmetric simple current in the first factor of the 3^5 model:

Note: Roughly said a simple current produces a new partition function thus new CFT from an given one.

Result:

- One supercharge from the left-movers, none from the right-movers \rightarrow L-R asymmetry, $\mathcal{N} = 1$ **target space SUSY**.
- 4 minimal factors unaffected \Rightarrow still **4 variables of weight 1**.
- Simple current splits first factor in **2 variables of weight 2**.

\Rightarrow **The model has still the structure of a $WC\mathbb{P}$ with $w_i = 1, 1, 1, 1, 2, 2$ and polynomials of degree 5!**

Educated guess: Is this the CFT to the $\mathcal{N} = 2$ SUGRA of IIB on $\mathbb{P}_{1,1,1,1,2,2}[5, 3]$ with SUSY breaking fluxes?

$\mathcal{N} = 2 \rightarrow \mathcal{N} = 1$ **breaking**: [Louis, Smyth, Triendl '09,10; Louis, Hansen '13]

- Needs **simultaneous geometric + non-geometric gaugings thus DFT**

No surprise: Our model is L/R asymmetric

- Resulting $\mathcal{N} = 1$ **spectrum is highly constrained**. For the above $P_{1,1,1,1,2,2}[5, 3]^{h_{12}, h_{11}=83, 2}$ only 6 possibilities:

$$(N_V, N_{\text{ax}}) \in \{(80, 0), (80, 1), (81, 0), (81, 1), (82, 1), (82, 2)\}$$

Compare: Our model has $(N_V; N_{\text{ax}}) = (80, 0) \checkmark$

Observation:

This ACFT looks like the (fully backreacted) string uplift of the GSUGRA of IIB on $P_{1,1,1,1,2,2}[5, 3] + (\text{SUSY breaking})$ fluxes!

More examples in our paper.

ACFT/GSUGRA conjecture:

A certain class of asymmetric Gepner models can be identified with the fully backreacted $\mathcal{N} = 1$ vacua of a $\mathcal{N} = 2$ GSUGRA.

Comments:

- We can compare the ACFT only to the *kinematics* of the GSUGRA, therefore its massless multiplet structure.
- Recall the flux scaling scenario from our group:
Non-geometric (thus winding) fluxes generically have a backreaction $\mathcal{O}(1)$ onto the geometry and no dilute flux limit ("want so shrink their cycle").

⇒ [Blumenhagen, Font, MF, Herschmann, Plauschinn, Sekiguchi, Wolf '15]

Our claim: After adjusting accordingly the minima of the GSUGRA can be uplifted to a full string solution.

ACFT/GSUGRA conjecture:

A certain class of asymmetric Gepner models can be identified with the fully backreacted $\mathcal{N} = 1$ vacua of a $\mathcal{N} = 2$ GSUGRA.

Consequences:

- Partial **SUSY breaking possible** beyond leading order.
- **Minima of GSUGRA** can correspond to **classical vacua of string theory**. The GSUGRA correctly captures the *kinematics* but is unlikely to describe the *dynamics* in a LEEA.
- **Non-geometric fluxes/gaugings (DFT!)** are part of the string dofs and **correspond to ACFTs**.

See also [Dabholkar, Hull '02,05; Condeescu, Florakis, Kounnas, Lüst '12,13]

Similar spirit: [Garcia-Etxebarria, Regalado]

Extended SUSY

Advantage: No superpotential, masses only through Higgs.
⇒ Perfect to test the conjecture in a more controlled setup

Ralphs talk: Asymmetric Gepner models in **6D & 8D**.

- $\mathbb{T}^2, K3$: **No NS-NS fluxes** supportable
- \mathbb{T}^4 : **SUSY breaking** kinematically **forbidden**

⇒ No explanation in terms of a GSUGRA possible!

✓ All models we found are (asymmetric) orbifolds of $\mathbb{T}^2, \mathbb{T}^4, K3$.

Note: $(-1)^{F_L}$ factor appeared often. E.g. the one from [Hellerman, McGreevy, Williams]

In the following: The more interesting 4D case!

4D: SUSY breaking

Still: No superpotential, only Higgs.

But: NS-NS fluxes supportable, more SUSY breakings allowed!

[Deser, Zumino '77; Cremmer, e.a. '78,79; Andrianopoli, D'Auria, Ferrara, Lledo '02]

Example: IIB on \mathbb{T}^6 has $\mathcal{N} = 8$ and therefore only the SUGRA multiplet $\mathcal{G}_{(8)}$ at the massless level. Some possible breakings:

$$\mathcal{G}_{(8)} \rightarrow \mathcal{G}_{(6)} + 2 \cdot \overline{\mathcal{S}}_{(6)}^{\frac{3}{2}\text{massive}}$$

$$\mathcal{G}_{(8)} \rightarrow \mathcal{G}_{(4)} + (6 - 2k) \cdot \mathcal{V}_{(4)} + k \cdot \overline{\mathcal{V}}_{(4)}^{\text{massive}} \quad k \in \mathbb{N}_0$$

Task

Run a stochastic search with up to 4 simultaneous simple currents to classify all asymmetric (non-geometric) Gepner models and see whether they are compatible with ACFT/GSUGRA! $\mathcal{O}(10^8)$ models!

Classification scheme:

$${}^D\mathfrak{N}_{[\mathcal{N}_L, \mathcal{N}_R]}$$

D : uncompactified directions

$\mathcal{N}_{L/R}$: counts the SUSY generators
from the L/R of the CFT.

Examples with only the SUGRA multiplet:

- ${}^4\mathfrak{N}_{[4,4]}$ thus $\mathcal{N} = 8$. The \mathbb{T}^6 compactification.
- ${}^4\mathfrak{N}_{[2,4]}$ thus $\mathcal{N} = 6$. Either broken $\mathcal{N} = 8$ or $\mathbb{T}^6/(\mathbb{Z}_2^L S)$
- ${}^4\mathfrak{N}_{[1,4]}$ thus $\mathcal{N} = 5$. Only interpretation is $\mathbb{T}^6/(\mathbb{Z}_2^L S, \tilde{\mathbb{Z}}_2^L \tilde{S})$

${}^4\mathfrak{N}_{[0,4]}$ has $\mathcal{N} = 4$

Massless spectrum:

$$\mathcal{G}_{(4)} + n_V \times \mathcal{V}_{(4)}, \quad n_V = 0, 2, 4, 6, 8, 10, 14, 18$$

$n_V = 18$: Asymmetric orbifold with $SU(2)^6$ gauge group

[Dixon, Kaplunovskiy, Vafa '87] $n_V = 6, \dots, 18$ Coloumb branch

And the rest?

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And the rest? Recall the super Higgs!

$$\mathcal{G}_{(8)} \rightarrow \mathcal{G}_{(4)} + (6 - 2k) \cdot \mathcal{V}_{(4)} + k \cdot \overline{\mathcal{V}}_{(4)}^{\text{massive}} \quad k \in \mathbb{N}_0$$

✓ perfect **match between GSUGRA and ACFT!**

Note: A priori no reason for steps of two in the CFT!

$${}^4\mathfrak{N}_{[1,2]} \text{ has } \mathcal{N} = 3$$

Massless spectrum:

$$\mathcal{G}_{(3)} + n_V \times \mathcal{V}_{(3)}, \quad n_V = 3, 7, 11, 13, 19$$

Fully explainable by super Higgs ($k \in \mathbb{N}_0$):

\mathcal{N}'	\mathcal{N}	massless spectrum
8 thus \mathbb{T}^6	3	$\mathcal{G}_3 + (3 - 2k) \cdot \mathcal{V}_3$
6	3	—
5	3	—
4 thus $\mathbb{T}^2 \times K3$	3	$\mathcal{G}_3 + (19 - 2k) \cdot \mathcal{V}_3$

$${}^4\mathfrak{N}_{[0,2]} \text{ has } \mathcal{N} = 2$$

Massless spectrum: $\mathcal{G}_{(2)} + n_V \times \mathcal{V}_{(2)} + n_H \times \mathcal{H}_{(2)}$

1. $n_H - n_V = 1$ with $n_V = 1, 3, 5, 6, \dots, 15, 17, \mathbf{19}, 20, 21, 22, 23$
2. $n_H - n_V = 13$ with $n_V = 3, 4, 5, \mathbf{7}, 8, 9, 10, 11$
3. $n_V - n_H = 11$ with $n_V = 13, 15, 17, \mathbf{19}, 21, 23$

Mechanisms at work:

- a) One basic model (**fat**) + gauge enhancement with up to 4 (higgsable) $\mathcal{V}_{(2)} + \mathcal{H}_{(2)}$ pairs (same as in 6D)
- b) SUSY breaking of $\mathbb{T}^2 \times K3$ thus $\mathcal{N} = 4 \rightarrow \mathcal{N} = 2$ with different number of short/long massive gravitino/vector multiplets. *Note the steps of one!*

1. $n_H - n_V = 1$ with $n_V = 1, 3, 5, 6, \dots, 15, 17, \mathbf{19}, 20, 21, 22, 23$

$n_V = 19$ is the \mathbb{T}^2 reduction of $\mathbb{T}^4 / \{\Theta, \Theta S(-1)^{F_L}\} \in {}^6\mathfrak{M}_{[1,0]}$

[Hellerman, McGreevy, Williams '04]; Θ reflection, S momentum shift

Again: Up to 4 additional $\mathcal{V} + \mathcal{H}$ pairs!

Alternatively $\mathcal{N} = 4 \rightarrow \mathcal{N} = 2$ breaking of $K3 \times \mathbb{T}^2$ with (only) short massive gravitino multiplets yields

$$n_V = 19 - k \qquad n_H = 20 - k \quad \checkmark$$

2. $n_H - n_V = 13$ with $n_V = 3, 4, 5, \mathbf{7}, 8, 9, 10, 11$

$\mathcal{N} = 4 \rightarrow \mathcal{N} = 2$ breaking of $K3 \times \mathbb{T}^2$ with no short gravitino and six short vector multiplets gives

$$n_V = 7 - k \qquad n_H = 20 - k \quad \checkmark$$

Alternatively, the model with $n_V = 7$ is the $K3$ reduction of the $\mathbb{T}^2/\{(-1)^{F_L}SW\} \in {}^8\mathfrak{N}_{[1,0]}$ model, therefore

$$\frac{\mathbb{T}^4 \times \mathbb{T}^2}{\{\mathbb{Z}_2, (-1)^{F_L}SW\}},$$

Allows for discrete torsion $\epsilon = \pm 1$ between the \mathbb{Z}_2 factors!
 $\epsilon = -1$ gives $n_V = 19$ and $n_H = 8$. Indeed:

3. $n_V - n_H = 11$ with $n_V = 13, 15, 17, \mathbf{19}, 21, 23$

Again in both cases: Up to 4 additional $\mathcal{V} + \mathcal{H}$ pairs!

For completeness: ${}^4\mathfrak{N}_{[2,2]}$ (symmetric) has $\mathcal{N} = 4$

Massless spectrum:

$$\mathcal{G}_{(4)} + (2 + n) \times \mathcal{V}_{(4)}, \quad n_V = 2 + n = 22, 14, 10, 6, 4$$

Clearly $n_V = 22$ is IIB on $\mathbb{T}^2 \times K3$. Rest is:

$$\text{Orb}_{n,m} = \frac{\mathbb{T}^4 \times \mathbb{T}^2}{\mathbb{Z}_n \mathcal{S}_m}$$

$\text{Orb}_{n,m}$	twisted sector vectors	massless spectrum
(2, 2)	$(1, \theta) = (6, 0)$	$\mathcal{G}_{(4)} + 6 \cdot \mathcal{V}_{(4)}$
(3, 3)	$(1, \theta, \theta^2) = (4, 0, 0)$	$\mathcal{G}_{(4)} + 4 \cdot \mathcal{V}_{(4)}$
(4, 2)	$(1, \theta, \theta^2, \theta^3) = (4, 0, 10, 0)$	$\mathcal{G}_{(4)} + 14 \cdot \mathcal{V}_{(4)}$
(6, 3)	$(1, \theta, \theta^2, \theta^3, \theta^4, \theta^5) = (4, 0, 0, 6, 0, 0)$	$\mathcal{G}_{(4)} + 10 \cdot \mathcal{V}_{(4)}$

Conclusion

- Concrete examples in 4D with $\mathcal{N} = 1$
- Classification of all asymmetric Gepner models in 4D, 6D and 8D with more SUSY

All examples support: **ACFT** \sim **GSUGRA**

Concretely: The asymmetric Gepner models we constructed correspond to fully backreacted minima of GSUGRA with geometric + non-geometric gaugings/fluxes.