

# Heterotic Calabi-Yau Flux Compactifications

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based on work together with Andre Lukas and Eirik Svanes

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compactifications and  
flux

Calabi-Yau domain walls

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# Motivation: Why flux on Calabi-Yau?

# Why heterotic Calabi-Yau?

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# Why heterotic Calabi-Yau?

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Compactification of the form

$$\mathcal{M}_{10} = \underbrace{\mathcal{M}_4}_{\text{maximally symmetric}} \times \underbrace{\mathcal{M}_6}_{\text{Calabi-Yau}}$$

# Why heterotic Calabi-Yau?

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## Compactification of the form

$$\mathcal{M}_{10} = \underbrace{\mathcal{M}_4}_{\text{maximally symmetric}} \times \underbrace{\mathcal{M}_6}_{\text{Calabi-Yau}}$$

- phenomenologically very attractive, recently several hundred standard models (spectrum and more) discovered, see [Anderson, Gray, Lukas, Palti; 1106.4804]
- powerful tools of algebraic geometry available

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Models attractive, but: moduli stabilisation is a problem

- in type IIA/B fluxes can help to resolve this
- heterotic has NS flux available (NS5 branes)

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Problem: NS flux on Calabi-Yau is very constrained (no-go, see below)



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Problem: NS flux on Calabi-Yau is very constrained (no-go, see below)

These approaches have been tried in the past:

- geometric flux [Gurrieri, Lukas, Micu 0408121], but:  $\mathcal{M}_6$  no longer Calabi-Yau)
- recent progress on this in standard CY compactifications without flux [Anderson, Gray, Lukas, Ovrut]
- ▶ this talk: Calabi-Yau *and* flux

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# Calabi-Yau compactifications and flux

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Make a compactification ansatz:

$$\mathcal{M}_{10} = \underbrace{\mathcal{M}_4}_{\text{maximally symmetric}} \times \mathcal{M}_6$$

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Make a compactification ansatz:

$$\mathcal{M}_{10} = \underbrace{\mathcal{M}_4}_{\text{maximally symmetric}} \times \mathcal{M}_6$$

then requiring unbroken supersymmetry gives [Wit, Smit, Dass, Nucl.Phys B283 1987]

$$H = 0 \quad \Leftrightarrow \quad \mathcal{M}_6 \text{ is Calabi-Yau}$$

This means either we give up Calabi-Yau or we have no flux.

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**?** Why is there such a tension between Calabi-Yau and flux?  
simple argument to illuminate this [Gauntlett, Martelli, Waldram 0302158]:  
look at dilaton equation of motion

$$\nabla^2 e^{-2\phi} = e^{-2\phi} * (H \wedge *H)$$

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$$\nabla^2 e^{-2\phi} = e^{-2\phi} * (H \wedge *H)$$

integrating gives

$$\underbrace{- \int_{X_6} d_6 \left( e^{4A} *_6 d_6 e^{-2\phi} \right)}_{=0} = \int_{X_6} e^{4A} e^{-2\phi} (H \wedge *_6 H) = \|e^{2A} e^{-\phi} H\|^2$$

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$$\Rightarrow H = 0 \quad \forall \mathcal{M}_6 \text{ compact}$$

i.e. we did not use Calabi-Yau, nor supersymmetry! (of course,  $\alpha'$  corrections avoid this argument, however: again no CY! [Strominger])

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We saw that the no-go result can be understood only making use of:

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Let us try relaxing this:

$$\mathcal{M}_{10} = \underbrace{\mathcal{M}_3}_{\text{maximally symmetric}} \times \mathcal{M}_7 =$$

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Now the dilaton equation of motion gives after integrating

$$-\partial_y^2 e^{-2\phi} - \partial_y e^{-2\phi} \frac{\partial_y V}{V} = \frac{1}{V} \|e^{-\phi} H\|^2$$

→ giving  $\phi$  an appropriate  $y$  dependence allows for non-zero flux!

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We have motivated the existence of Calabi-Yau domain walls with flux. Let us make this more explicit:

$$\left(\nabla_M + \frac{1}{8}\mathcal{H}_M\right)\epsilon = 0 \quad \left(\nabla\phi + \frac{1}{12}\mathcal{H}\right)\epsilon = 0$$

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$$\left(\nabla_M + \frac{1}{8}\mathcal{H}_M\right)\epsilon = 0 \quad \left(\nabla\phi + \frac{1}{12}\mathcal{H}\right)\epsilon = 0$$

using  $6d$  spinors we define the forms on  $\mathcal{M}_6$ :

$$J := \eta_-^\dagger \gamma_{uv} \eta_- e^{uv} \quad \Omega := \eta_+^\dagger \gamma_{uvw} \eta_- e^{uvw}$$

with the domain wall ansatz  $\mathcal{M}_{10} = \mathcal{M}_3 \times \mathbb{R} \times \mathcal{M}_6$  this leads to

$$\begin{aligned} d\Omega_- &= 2d\phi \wedge \Omega_- & J \wedge dJ &= J \wedge J \wedge d\phi \\ J \wedge H &= *d\phi & dJ &= 2\phi' \Omega_- - \Omega'_- - 2d\phi \wedge J + *H \\ d\Omega_+ &= J \wedge J' - \phi' J \wedge J + 2d\phi \wedge \Omega_+ & \Omega_- \wedge H &= 2\phi' *1 \\ \Omega_+ \wedge H &= 0, \end{aligned}$$

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The Killing spinor equations then reduce for a Calabi-Yau (i.e.  $dJ = 0$ ,  $d\Omega = 0$ ) to the following flow equations

$$\begin{aligned}\Omega'_+ &= 2\phi'\Omega_+ - H \\ J \wedge J' &= \phi' J \wedge J \\ \Omega_- \wedge H &= 2\phi' * 1 ,\end{aligned}$$

together with the constraint

$$\Omega_+ \wedge H = 0 .$$

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→ existence of solutions guaranteed by [Hitchin math.DG 0107101]



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Expand forms into a basis of harmonic three- and two-forms,  $\{\alpha_A, \beta^B\}$  and  $\omega_i$

$$\Omega = Z^A \left[ \alpha_A - \mathcal{G}_{AB}(Z^A) \beta^B \right]$$

$$J = v^i \omega_i$$

$$H = \mu^A \alpha_A + \epsilon_B \beta^B$$

Here  $Z^A$  are the  $h^{2,1} + 1$  projective complex moduli and  $\mathcal{G}$  their prepotential,  $v^i$  the  $h^{1,1}$  Kähler moduli and  $\mu^A, \epsilon_A$  y-independent flux parameters.

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The flow equations for the forms can be easily integrated if we define  $X^A := e^{-2\phi} Z^A$  and a new coordinate  $z$  via  $dy/dz = e^{2\phi}$  :

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$$\operatorname{Re} X^A = -\mu^A z - \gamma^A$$

$$\operatorname{Re} \mathcal{G}_B(X^A) = \epsilon_B z + \eta_B$$

$$v^i = e^\phi v_0^i,$$

with integration constants  $\{\gamma^A, \eta_B\}$ , such that  $\gamma^A \epsilon_A + \eta_B \mu^B = 0$ .

$2(h^{2,1} + 1)$  equations for  $h^{2,1} + 1$  complex  $X^A \Rightarrow X^A = X^A(z)$ .

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→ dilaton can then be obtained by solving the ordinary first order differential equation

$$\operatorname{Im} X^A \epsilon_A + \operatorname{Im} \mathcal{G}_B \mu^B = -2V_0 \partial_z (e^{-\phi})$$

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- for any given Calabi-Yau such a solution can be constructed
- it allows for any harmonic flux to be present (no  $\alpha'$  corrections needed)

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- We saw that relaxing the  $4d$  spacetime to a domain wall gave us enough freedom to add flux to a given Calabi-Yau compactification but: did we lose all potential to do realistic phenomenology?

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to think further about this, let us look at the low-energy theory:



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$\mathcal{N} = 1$  SUGRA, moduli fields  $(S, T^i, X^A)$  [Gurrieri, Lukas, Micu 0408121]

$$K = -\ln i(\bar{S} - S) - \ln 8V - \ln i(X^A \bar{\mathcal{G}}_A - \bar{X}^A \mathcal{G}_A)$$

(so far: this is the same theory as for a maximally symmetric CY compactification)

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new: superpotential

$$W = \epsilon_A X^A + \mu^A \mathcal{G}_A$$

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The theory described above has for non-vanishing flux a  $1/2$ -BPS domain wall [Lukas, Matti 1005.5302].

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→ for the more general case half-flat: Yes, in the large complex structure limit [Lukas, Matti]

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Matching requires field redefinitions:

$$e^{2\phi} = e^{2\phi_4} V/V_0 \quad Z^A = e^{2\phi} X^A \quad v^i = t^i$$

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We also know something about the asymptotics:

weak coupling limit as  $y \rightarrow \infty$

$X^A$  approach constant value as  $y \rightarrow \infty$

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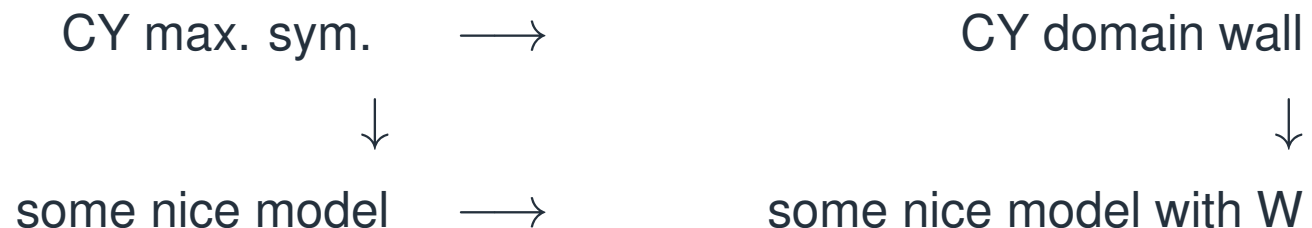
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Therefore, want to take the viewpoint:





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- domain wall is a perturbative solution and unstable: at least dilaton is still unfixed

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- domain wall is a perturbative solution and unstable: at least dilaton is still unfixed
- in standard Calabi-Yau compactification the situation is similar: at least dilaton unfixed

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- domain wall is a perturbative solution and unstable: at least dilaton is still unfixed
  - in standard Calabi-Yau compactification the situation is similar: at least dilaton unfixed
- add non-perturbative effects to lift the ground state to a stable vacuum

# domain wall = end of phenomenology?

## Motivation

Calabi-Yau compactifications and flux

## Calabi-Yau domain walls

Low energy phenomenology

$4d$  effective theory

domain wall ground states

domain wall = end of phenomenology?

Conclusions and Outlook

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However, why shouldn't the domain wall itself be lifted in this process, leading to a vacuum with a maximally symmetric spacetime in  $4d$ ?

In fact, for half-flat domain walls it was shown that the domain wall can indeed be lifted to a maximally symmetric and stable vacuum (moduli at consistent values) [M.K., Lukas, Matti, Svanes 1210.5933]

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# Conclusions and Outlook

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## Conclusions:

- heterotic CY compactifications offer a fertile ground for phenomenology
- but: moduli stabilisation problematic



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- heterotic CY compactifications offer a fertile ground for phenomenology
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- in particular CY domain walls allow for arbitrary harmonic flux

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## Conclusions:

- heterotic CY compactifications offer a fertile ground for phenomenology
- but: moduli stabilisation problematic
- relaxing assumptions on  $4d$  spacetime allows for flux
- in particular CY domain walls allow for arbitrary harmonic flux
- this results in a model which has a domain wall, so it seems no realistic phenomenology possible
- however, we believe that past work has shown that there is justified reason to believe that such a model can be lifted to a maximally symmetric vacuum

# Outlook: End of a paradigm?

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Paradigm:



make compactification ansatz as close as possible to our universe (in particular  $4d$  spacetime should be maximally symmetric)

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However:

- it seems very hard to actually find stable vacua with this approach
- it almost seems unavoidable to introduce non-perturbative effects to stabilise the dilaton

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However:

- it seems very hard to actually find stable vacua with this approach
- it almost seems unavoidable to introduce non-perturbative effects to stabilise the dilaton

What if we take a different point of view:

!?! start with some ansatz that allows for enough freedom to stabilise all moduli but the dilaton (= switch on a superpotential)

# Outlook: End of a paradigm?

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Paradigm:

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However:

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- it almost seems unavoidable to introduce non-perturbative effects to stabilise the dilaton

What if we take a different point of view:

- !?
- lift to non-perturbative stable vacuum, with a maximally symmetric spacetime

After all, it is the final lifted vacuum which we want to look like our universe.

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Still many things left to do:

- study lifting and moduli stabilisation for an explicit model
- try supersymmetric cosmic string, black hole, ...

# Thank you!

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Thank you very much!