Quartum state reconstruction with RBMS Consider first a Classical Ising model:  $H = -\sum_{z} \sigma_{z}^{z} \sigma_{z}^{z}$ With Boltzmann  $P(\vec{\sigma}) = \pm e^{-\beta E(\vec{\sigma})}$ Makes sense that an RBM should be able to veconstruct this. Positive real wavefunctions Imagine some quantum system where a basis can be chosen such that  $|4\rangle = \sum_{\sigma} |\sigma \times \sigma| |4\rangle = \sum_{\sigma} |4\rangle |\sigma\rangle$ and all 4(0) EIR and 4(0) >0 The Bon rule states P(o) = 14(o) 2 Thus we can construct an RBM such that

Care 2: d6-diagonal observables.  $\langle \hat{\mathcal{O}}^{oo} \rangle = \langle \mathcal{U}_{\lambda} | \hat{\mathcal{O}}_{oo'} | \mathcal{U}_{\lambda} \rangle$ = ZZ4, (0') (0') (0') (0') (10) 4, (0)  $= \sum_{n} \sum_{n} \int_{\lambda} (\sigma) \int_{\lambda} (\sigma)$  $= \sum_{\sigma} P_{\lambda}(\sigma) \sum_{\sigma'} \frac{\int P_{\lambda}(\sigma')}{\int P_{\lambda}(\sigma)} \mathcal{O}_{\sigma\sigma'}$ "local" estimate of O As long the Oron matrix is sufficiently sparse it's expectation value can be estimated by MCMC. You can see here how sufficient generalization is required. Q: For which Hamiltonicens are wavefunctions always the real?

Complex Wavefunctions More generally  $4(\vec{\sigma}) = |4(\vec{\sigma})| e^{i\phi(\vec{\sigma})}$ Clearly more bases are needed - P(3)=14(3)12 does not contain fingerprints of the phase of. G: How many bases in general? Typically? How to parameterize a complex wavefunction on on RBM? - any function approximate (FFNN, CNN) - Complex weights (Carleo) - another RBM (say with parameters µ) e.g.  $\mathcal{U}_{\lambda\mu}(\vec{\sigma}) = \int p(\vec{\sigma}) e^{i\phi_{\mu}(\vec{\sigma})} , \phi_{\mu}(\vec{\sigma}) = \log p_{\mu}(\vec{\sigma})$ How to train? 'Target "(unknown) weighterent must produce data in different bases (unitary transformation)  $2(\overline{c}^{b}) = \sum_{\sigma} 2(\overline{c}^{b}, \overline{c}, \overline{c}, \overline{c})$ 

where 
$$|\vec{\partial}b\rangle = |\vec{\sigma}_{1}^{b_{1}}, \vec{\sigma}_{2}^{b_{0}}, \dots, \vec{\sigma}_{p}^{b_{n}}\rangle$$
  
ie each spin can be in a different boos.  
Phase information is then transmitted via  
 $P(\vec{\sigma}b) = |2f(\vec{\partial}b)|^{2}$ 

Then given a data set 
$$\mathcal{P} = \{\vec{\sigma}^{b}\}$$
 we can  
use a sum of KL divergences in different  
bases  
 $C_{A\mu} = -\frac{1}{||\mathcal{P}||} \sum_{\vec{\sigma}^{b} \in \mathcal{P}} \log |\mathcal{V}_{A\mu}(\vec{\sigma}^{b})|^{2}$   
 $= -\frac{1}{||\mathcal{P}||} \sum_{\vec{\sigma}^{b} \in \mathcal{P}} \left[\log (2\ell_{\vec{\sigma}^{b}}, \vec{\sigma}^{b}, \vec{\tau}_{A\mu}(\sigma)) + c.c.\right]$   
etc.

Geneally, it is not known how many basis are needed (maximum 2<sup>N</sup>), or how many measurements per basis are required...