



**T II: Elektrodynamik**  
(Prof. E. Frey)

**Problem set 11**

**Tutorial 11.1** *Hertz potential*

- a) Starting from Maxwell's equations introduce electromagnetic potentials and derive the wave equations

$$\left(\frac{1}{c^2}\partial_t^2 - \nabla^2\right)\vec{A} = \frac{4\pi}{c}\vec{j}, \quad \left(\frac{1}{c^2}\partial_t^2 - \nabla^2\right)\varphi = 4\pi\rho$$

provided the potentials fulfill the Lorentz gauge  $c^{-1}\partial_t\varphi + \text{div}\vec{A} = 0$ . Hence, the fields appear to be decoupled, but the gauge constraint has to be satisfied.

- b) The continuity equation for the charge is solved automatically by introducing the polarization  $\vec{P}$  with  $\partial_t\vec{P} = \vec{j}$  and  $-\text{div}\vec{P} = \rho$ . Show that the Lorentz gauge is satisfied automatically by introducing the (electric) Hertz potential  $\vec{Z}$  with  $c^{-1}\partial_t\vec{Z} = \vec{A}$  and  $-\text{div}\vec{Z} = \varphi$ . What equation governs the dynamics of  $\vec{Z}$ ? Express the electric and magnetic field in terms of the Hertz potential.
- c) Perform a temporal Fourier transform and show that the solution compatible with Sommerfeld's radiation condition is given by

$$\vec{Z}_\omega(\vec{x}) = \int d^3\vec{y} \frac{\vec{P}_\omega(\vec{y})}{|\vec{x} - \vec{y}|} e^{ik|\vec{x} - \vec{y}|}, \quad k = \frac{\omega}{c}.$$

- d) Argue that the leading behavior of the Hertz potential far away from the source is obtained by approximating

$$\vec{Z}_\omega(\vec{x}) = \frac{e^{ikr}}{r} \vec{g}(k\hat{n}), \quad \vec{g}(\vec{k}) = \int d^3\vec{y} e^{-i\vec{k}\cdot\vec{y}} \vec{P}_\omega(\vec{y}),$$

where  $\vec{x} = r\hat{n}$ ,  $|\hat{n}| = 1$  and determine the electric and magnetic field in the radiation zone.

- e) In the case of a highly localized source,  $kd \ll 1$ , expand the exponential in  $\vec{g}(\vec{k})$  in powers of  $\vec{k}$  and show that the leading contributions to the angular dependence of the radiation are of the following form,

$$\vec{g}(\vec{k}) = \vec{p} - \hat{n} \times \vec{m} - \frac{1}{6}i\vec{k} \cdot \underline{\underline{Q}} \dots$$

Interpret the individual terms.

**Problem 11.2**     *Magnetic Hertz potential*

In addition to the electric Hertz potential  $\vec{Z}$  one may also introduce a magnetic Hertz potential  $\vec{Z}_m$ .

- a) Show that the Lorentz gauge is automatically fulfilled provided one chooses  $\varphi = -\text{div } \vec{Z}$  and  $\vec{A} = c^{-1} \partial_t \vec{Z} + \text{curl } \vec{Z}_m$ . Show that the wave equations for the gauge potentials  $\varphi, \vec{A}$  are fulfilled if one requires

$$\left( \frac{1}{c^2} \partial_t^2 - \nabla^2 \right) \vec{Z} = 4\pi \vec{P} \quad \text{and} \quad \left( \frac{1}{c^2} \partial_t^2 - \nabla^2 \right) \vec{Z}_m = 4\pi \vec{M}.$$

Express the electromagnetic fields in terms of the two Hertz potentials.

- b) Perform a temporal Fourier transform of the wave equation for the magnetic Hertz potential and give a formal solution that fulfills Sommerfeld's radiation condition. Show that the leading behavior of the magnetic Hertz potential in a long-wavelength expansion far away from the source is given by

$$Z_m^\omega(\vec{x}) = \frac{e^{ikr}}{r} \vec{m}_\omega, \quad r = |\vec{x}|, \quad k = \omega/c,$$

and interpret the vector  $\vec{m}_\omega$ .

- c) Derive the corresponding electromagnetic fields  $\vec{E}, \vec{B}$  as well as the time-averaged Poynting vector in the radiation zone. Discuss the angular dependence of the emitted radiation and the polarization of the electric field for the case of  $\vec{m}_\omega = m_\omega(0, 0, 1)$ .

**Problem 11.3**     *Gravitational waves*

As a relativistic invariant minimal generalization of Newton's theory of gravity consider the schematic model

$$\left[ \frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \nabla^2 \right] \Phi(\vec{x}, t) = 4\pi G \rho(\vec{x}, t),$$

where  $\Phi(\vec{x}, t)$  denotes the gravitational potential,  $\rho(\vec{x}, t)$  the mass density,  $G$  the universal gravitational constant, and  $c$  the speed of light.

- a) Perform a temporal Fourier transform,  $\Phi_\omega(\vec{r}) = \int dt e^{i\omega t} \Phi(\vec{r}, t)$ , and show that the (causal) solution for a localized source is given by

$$\Phi_\omega(\vec{x}) = G \int d^3y \frac{\rho_\omega(\vec{y})}{|\vec{x} - \vec{y}|} e^{ik|\vec{x} - \vec{y}|}, \quad k = \omega/c.$$

- b) Argue that the leading behavior of the potential in the radiation zone  $r \gg \lambda$  and far away from the localized source  $r \gg d$  is obtained by approximating

$$\Phi_\omega(\vec{x}) \simeq G \frac{e^{ikr}}{r} \int d^3y e^{-ik\hat{n} \cdot \vec{y}} \rho_\omega(\vec{y}), \quad \vec{x} = r\hat{n}, \quad |\hat{n}| = 1.$$

- c) If the linear dimensions  $d$  of the radiating source are small compared to the wavelength,  $kd \ll 1$ , the exponential  $e^{-ik\hat{n} \cdot \vec{y}}$  can be expanded in a power series. Determine the lowest order contribution in this long-wavelength expansion reflecting that the mass and momentum conservation laws hold

$$\partial_t \rho(\vec{x}, t) + \nabla_i J_i(\vec{x}, t) = 0, \quad \partial_t J_i(\vec{x}, t) + \nabla_j \Pi_{ij}(\vec{x}, t) = 0,$$

where  $\vec{J}(\vec{x}, t)$  denotes the mass current also identified with the momentum density, and  $\Pi_{ij}(\vec{x}, t)$  corresponds to the momentum current.

**Problem 11.4**    *Rotating dipole*

Consider an electric dipole rotating in the  $x$ - $y$  plane,

$$\vec{p}(t) = \text{Re} (\vec{p}_\omega e^{-i\omega t}) , \quad \vec{p}_\omega = p_\omega(1, i, 0)/\sqrt{2} ,$$

and discuss the emitted electromagnetic radiation.

- a) Determine the angular dependence of the time-averaged Poynting vector for radiation emitted in the direction  $\hat{n} = (\sin \vartheta \cos \phi, \sin \vartheta \sin \phi, \cos \vartheta)$ . What is the total power radiated by the dipole?
- b) Evaluate the polarization of the emitted radiation for directions  $\hat{n}$ 
  - (i) in the plane of rotation ( $\vartheta = \pi/2$ ),
  - (ii) perpendicular to it ( $\vartheta = 0$ ).

*Due date: Tuesday, 7/10/2007, at 9 p.m.*