

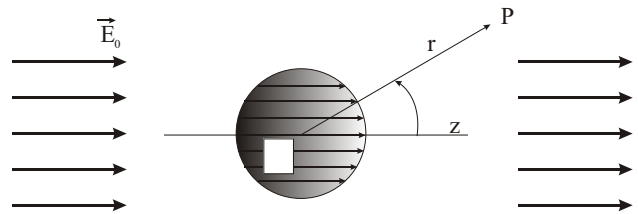


T II: Elektrodynamik
(Prof. E. Frey)

Problem set 10

Tutorial 10.1 Dielectric sphere

A dielectric sphere of radius R characterized by a dielectric constant ϵ is placed in an initially uniform electric field \vec{E}_0 . For convenience, choose the center of the sphere as the origin and consider \vec{E}_0 along the z -axis. Then the problem exhibits an axial symmetry, which simplifies the problem.



- a) Determine the electrostatic potential $\varphi(\vec{x})$ inside the sphere, $|\vec{x}| \leq R$, and outside the sphere, $|\vec{x}| > R$. Formulate appropriate matching conditions at the surface of the sphere. Recall that the most general axially symmetric solution of Laplace's equation $\nabla^2 \varphi = 0$ in polar coordinates, $\vec{x} = r(\sin \vartheta \cos \phi, \sin \vartheta \sin \phi, \cos \vartheta)$, is given in terms of

$$\varphi(\vec{x}) = \sum_{\ell=0}^{\infty} \left(a_{\ell} r^{\ell} + b_{\ell} r^{-(\ell+1)} \right) P_{\ell}(\cos \vartheta),$$

where $P_{\ell}(t)$ denotes the Legendre polynomials and a_{ℓ}, b_{ℓ} are undetermined coefficients.

- b) Derive the corresponding electric field $\vec{E}(\vec{x})$. Extract the polarization $\vec{P}(\vec{x})$ inside the sphere and find the total induced dipole moment \vec{p} of the dielectric sphere. Determine the effective polarizability of the sphere α defined by $\vec{p} = \alpha \vec{E}_0$.
- c) Show that charges accumulate at the surface of the sphere and determine the induced surface charge density σ .

Problem 10.2 *Magnetic shielding – μ -metal*

A μ -metal is a nickel-iron alloy that has a very high magnetic permeability $\mu \sim 10^4 - 10^6 \gg 1$. The technical application of these materials is the screening of static (or low-frequency) magnetic fields, which cannot be attenuated by other methods.

Consider a spherical shell of magnetic permeability μ and inner and outer radii R_i and R_o , respectively, placed in a previously uniform magnetic field \vec{H}_∞ . The medium inside and outside of the shell has a magnetic permeability $\mu = 1$.

- Argue that one can introduce a scalar magnetic potential φ_M to represent the field $\vec{H} = -\vec{\nabla}\varphi_M$, and show that it fulfills the Laplace equation $\nabla^2\varphi_M = 0$ in each region.
- State the appropriate matching conditions for φ_M at the interfaces $r = R_i$ and $r = R_o$.
- Recalling that the most general solution of the Laplace equation with cylindrical symmetry is provided by

$$\varphi_M(\vec{r}) = \sum_{\ell=0}^{\infty} \left(a_\ell r^\ell + b_\ell r^{-(\ell+1)} \right) P_\ell(\cos \vartheta),$$

determine the magnetostatic potential in each region. Calculate explicitly the corresponding magnetic field \vec{H} inside of the shell.

- Determine the leading behavior of the field inside of a thin shell of a μ -metal for $\mu \rightarrow \infty$. Discuss why a μ -metal provides an effective shielding.

Problem 10.3 *Legendre polynomials*

Consider the following partial differential equation

$$\frac{\partial}{\partial t} \left[(1-t^2) \frac{\partial \psi}{\partial t} \right] = -\frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right). \quad (*)$$

A solution is provided by the function

$$\psi(r, t) = \frac{1}{\sqrt{1-2rt+r^2}}, \quad -1 < r < 1, \quad -1 \leq t \leq 1,$$

which serves as a generating function for the *Legendre polynomials* $P_\ell(t)$, i.e., a Taylor expansion with respect to r , $\psi(r, t) = \sum_{\ell=0}^{\infty} r^\ell P_\ell(t)$, defines the functions $P_\ell(t)$.

- Show by explicit substitution that $\psi(r, t)$ indeed solves the partial differential equation (*).
- Identifying $t = \cos \vartheta$ reveals that $\psi(r, t)$ corresponds to the Coulomb potential of a unit charge located on the z -axis at unit distance from the origin. Thus $\psi(r, t = \cos \vartheta)$ solves the Laplace equation in polar coordinates

$$\nabla^2 \psi(r, \cos \vartheta) = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2 \sin \vartheta} \frac{\partial}{\partial \vartheta} \left(\sin \vartheta \frac{\partial \psi}{\partial \vartheta} \right) = 0.$$

Using this observation derive the partial differential equation (*).

- Determine explicitly $P_\ell(t = \pm 1)$ observing that for $t = \pm 1$ the Taylor series of $\psi(r, t)$ becomes elementary.
- Employ the symmetries of $\psi(r, t)$ to argue that $P_\ell(t)$ is a symmetric (anti-symmetric) function for even (odd) ℓ .
- Inspect the Taylor series to demonstrate that $P_\ell(t)$ is a polynomial of order ℓ .
Hint: Expand the square root in $x = 2rt - r^2$.
- Calculate and sketch the first four Legendre polynomials ($\ell = 0, \dots, 3$).

- g) Substitute the Taylor series of $\psi(r, t)$ in the partial differential equation (*). Comparing the coefficients of r^ℓ confirm that the $P_\ell(t)$ satisfy the second order differential equation, i.e., they are indeed Legendre polynomials,

$$\frac{d}{dt} \left[(1-t^2) \frac{dP_\ell(t)}{dt} \right] + \ell(\ell+1)P_\ell(t) = 0, \quad -1 \leq t \leq 1. \quad (**)$$

- h) Show that $(1-2rt+r^2)\partial\psi/\partial r = (t-r)\psi$. Make use of this result to derive the recursion relation

$$\ell P_{\ell-1}(t) - (2\ell+1)tP_\ell(t) + (\ell+1)P_{\ell+1}(t) = 0.$$

Similarly, verify that $(1-2rt+r^2)\partial\psi/\partial t = r\psi$ and prove

$$P'_{\ell+1}(t) - 2tP'_\ell(t) + P'_{\ell-1}(t) = P_\ell(t).$$

- i*) Show that the Legendre polynomials are orthogonal in the following sense,

$$\int_{-1}^1 dt P_\ell(t)P_{\ell'}(t) = \frac{2}{2\ell+1}\delta_{\ell\ell'}.$$

Hint: Employ the differential equation (**) to show that

$$[\ell'(\ell'+1) - \ell(\ell+1)] \int_{-1}^1 dt P_\ell(t)P_{\ell'}(t) = 0$$

and conclude that orthogonality holds. The normalization follows by considering $\int dt \psi(r, t)^2$. First perform the integration directly; then use the Taylor expansion in r and the orthogonality property.

Due date: Tuesday, 7/3/07, at 9 p.m.